

# Energetics of $\text{H}_2$ clusters from density functional and coupled cluster theories

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We use coupled-cluster quantum chemical methods to calculate the energetics of molecular clusters cut out of periodic molecular hydrogen structures that model observed phases of solid hydrogen. The hydrogen structures are obtained from Kohn-Sham density functional theory (DFT) calculations at pressures of 150, 250 and 350 GPa, which are within the pressure range in which phases II, III and IV are found to be stable. The calculated deviations in the DFT energies from the coupled-cluster data are reported for different functionals, and optimized functionals are generated which provide reduced errors. We give recommendations for semi-local and hybrid density functionals that are expected to provide an accurate description of hydrogen at high pressures.

The phase diagram of solid hydrogen is of great theoretical and experimental interest. For example, very recent experiments by Dias and Silvera have claimed metallization of strongly compressed hydrogen, giving rise to world-wide attention and controversy<sup>1</sup>.

A number of theoretical approaches for exploring the structures and energetics of hydrogen have been used, such as density functional theory (DFT)<sup>2</sup>, diffusion quantum Monte Carlo (DMC)<sup>3,4</sup>, and path integral molecular dynamics (PIMD) methods<sup>5-8</sup>. The hexagonal structure of phase I of solid hydrogen was obtained from experiment long ago<sup>9,10</sup>. Structures for solid hydrogen at high-pressures have been proposed from the DFT-based *ab initio* random structure searching method<sup>11-13</sup>. These structures provide models for phases III<sup>14</sup>, IV<sup>15</sup> and, with less certainty, for phase II<sup>16</sup> of solid hydrogen which are in reasonable agreement with the available vibrational data from infra-red and Raman spectroscopy<sup>17-24</sup>. A hexagonal structure has recently been proposed as a model for the low-pressure range of phase III<sup>25</sup>, but it is not considered in this work. DMC has been used to benchmark DFT results for solid hydrogen at high pressures<sup>26-30</sup>, although the expense of the DMC calculations means that only very limited data could be generated.

Relative energies of hydrogen phases calculated using a range of density functionals exhibit variations that are considerably larger than the energy range of interest, which is  $\sim 10$  meV per proton. The large vibrational energies in solid hydrogen also lead to important corrections to the energetics, which are vital for interpreting the experimental data. In addition, X-ray and neutron diffraction experiments are very difficult at high pressures. In practise, only DFT methods can currently be used to perform calculations which provide a reasonable description of the phase diagram of hydrogen. It is therefore crucial to design density functionals for compressed hydrogen that are accurate at the high pressures reached in diamond-anvil-cell experiments.

Here we construct exchange-correlation functionals that can be expected to provide an accurate description of the energetics of the structures mentioned above within DFT. Two accurate hybrid functionals are provided, denoted  $\text{O}_1$  and  $\text{O}_1\text{-D3}$ , for use without and with

dispersion corrections, respectively. We also provide accurate generalized gradient approximation (GGA) functionals (though less accurate than their hybrid counterparts), denoted  $\text{O}_3$  and  $\text{O}_3\text{-D3}$ , for use without and with dispersion corrections, respectively. DMC energies are computed using orbitals generated with one of the functionals and compared with results from coupled cluster singles, doubles, and perturbative triples [CCSD(T)] in order to further assess the accuracy of the hydrogen functionals.

## I. GENERATION OF $\text{H}_2$ CLUSTER GEOMETRIES

We use small molecular hydrogen clusters as reference systems to construct our DFT functionals. We extract cluster geometries from stable bulk DFT structures obtained with the PBE semi-local density functional by taking spherical cuts of each system containing up to 24 hydrogen atoms. Total energy calculations for a variety of clusters of hydrogen molecules are carried out at the level of CCSD(T) and DFT. The coupled cluster and DFT calculations are performed using the MOLPRO code<sup>31</sup>.

The procedure for generating cluster geometries is designed to give reasonably compact and symmetric structures and is as follows. For a bulk structure, atomic pairs closer than  $0.85 \text{ \AA}$  are classified as molecules. The distance parameter used in this classification is chosen so that all atoms are classified as a member of only one molecule. The center of mass of each  $\text{H}_2$  molecule and the midpoint between the center of mass of pairs of molecules is then recorded.

For clusters with an even number of molecules, a sphere is centered at each midpoint and expanded until it contains the required number of hydrogen atoms,  $N$ . For clusters with an odd number of molecules, a sphere is placed at the center of mass of each  $\text{H}_2$  molecule and expanded until it contains  $N$  hydrogen atoms. This procedure is carried out for each midpoint/center-of-mass, and we select the point about which the magnitude of the first moment of the selected atomic positions is minimized. This process preserves much of the symmetry

of the underlying bulk structure, and supplies compact clusters representing each structure.

Clusters are generated from the bulk structures of symmetries  $P2_1c$ -24,  $C2/c$ -24, and  $Pc$ -48, which model, respectively, phases II, III, and IV of hydrogen, while  $Cmca$ -12 and  $Cmca$ -4 are other low-energy structures that have been found in structure searches. (Note that in this notation the number of atoms in the primitive unit cell is given following the space group designation.) The five structures are calculated using DFT at pressures of 150, 250, and 350 GPa, which covers the experimentally relevant high-pressure regime. Clusters containing 1–12 pairs of hydrogen atoms are generated from each of these structures, providing a total of 180 clusters.

## II. EXTRAPOLATION AND ERROR ASSESSMENT FOR CCSD(T) AND DFT

CCSD(T) calculations are performed to provide reference energies for each cluster which accurately include many-body effects. These reference energies are then compared with DFT energies for several exchange-correlation functionals. The cc-pVnZ Gaussian basis sets are used<sup>32</sup>, together with extrapolation to the complete basis set (CBS) limit. Note that  $n = D, T, Q, 5, 6, \dots$  specifies the largest value of the angular momentum index  $l$  in the basis set.

A variety of procedures are available in the literature to extrapolate CCSD(T) results to the CBS limit<sup>33</sup>. We employ a computationally expensive extrapolation procedure, which we refer to as “accurate extrapolation”, for small clusters containing 2–8 hydrogen atoms. These “accurate” total energies are then used to validate an empirical, but computationally cheaper, extrapolation procedure that is applied to all cluster sizes, which we refer to as “efficient extrapolation”.

The CCSD(T) energy is given by

$$E_{\text{CCSD(T)}} = E_{\text{HF}} + E_{\text{S}} + E_{\text{D}} + E_{\text{(T)}} , \quad (1)$$

where  $E_{\text{HF}}$  is the Hartree-Fock (HF) energy,  $E_{\text{S}}$  is the second-order contribution to the correlation energy due to double excitations of opposite-spin electrons,  $E_{\text{D}}$  is the second-order contribution to the correlation energy due to double excitations of like-spin electrons, and  $E_{\text{(T)}}$  is the perturbative triple-excitation correlation energy. We extrapolate the CCSD(T) energies of the smaller clusters to the CBS limit using analytic forms justified by the expected asymptotic behavior of the components of the CCSD(T) energy<sup>34</sup>,

$$E_{\text{HF}}(n) = E_{\text{HF}}(\infty) + a_0 e^{-b_0 \sqrt{n}} , \quad (2)$$

$$E_{\text{S}}(n) = E_{\text{S}}(\infty) + a_1 (n + c_1)^{-3} + b_1 (n + c_1)^{-5} , \quad (3)$$

$$E_{\text{D}}(n) = E_{\text{D}}(\infty) + a_2 (n + c_2)^{-5} + b_2 (n + c_2)^{-7} , \quad (4)$$

$$E_{\text{(T)}}(n) = E_{\text{(T)}}(\infty) + a_3 (n + c_3)^{-3} + b_3 (n + c_3)^{-5} , \quad (5)$$

where the left-hand sides of the equations are the components of the CCSD(T) energy evaluated with the cc-

pVnZ basis set,  $\{a_i, b_i, c_i\}$  are fitting parameters describing the variation of the energy components with  $n$ , and  $E_{\text{HF}}(\infty)$ ,  $E_{\text{S}}(\infty)$ ,  $E_{\text{D}}(\infty)$ , and  $E_{\text{(T)}}(\infty)$  are the extrapolated energy components, which are treated as fitting parameters. Values of the left-hand sides are collected from CCSD(T) calculations using basis sets with  $n = T, Q, 5, 6$ , and a least-squares fit for each component then provides the estimated CBS total energy.

This extrapolation is expected to be accurate due to the analytically justified forms for each component of the total energy, and because including higher order terms in the correlation energy expressions ensures that the variation of energy with basis set size is accurately described for the smaller basis sets. The absolute deviation  $|E_{\text{CCSD(T)}}(6) - E_{\text{CCSD(T)}}(\infty)|$  is expected to severely overestimate the error in this extrapolation process, hence the mean and maximum absolute deviations of 4.6 and 5.4 meV/[H<sub>2</sub>] for clusters of 2–8 hydrogen atoms reliably validate this finite basis correction.

Next, we use these “accurate” total energies to validate a computationally cheaper extrapolation method for application to the full set of clusters. The variation of the total energy with basis set is dominated by the  $E_{\text{S}}(n)$  term, hence we take the lowest-order part of  $E_{\text{S}}(n)$  to construct our estimated CBS limit,  $E_{\text{est}}$ . We also estimate an extrapolation error interval by using similar forms with different powers to define upper and lower limits,  $E_{\text{u}}$  and  $E_{\text{l}}$ , with values of the powers chosen so that the estimated range includes the “accurate” estimates for the small clusters of 2–8 hydrogen atoms.

The resulting “efficient extrapolation” procedure takes the form

$$E_{\text{u}}(n) = E_{\text{u}}(\infty) + a_{\text{u}} (n + n_0)^{-\frac{5}{2}} , \quad (6)$$

$$E_{\text{est}}(n) = E_{\text{est}}(\infty) + a_{\text{est}} (n + n_0)^{-3} , \quad (7)$$

$$E_{\text{l}}(n) = E_{\text{l}}(\infty) + a_{\text{l}} (n + n_0)^{-\frac{7}{2}} , \quad (8)$$

where  $E_{\text{u}}(\infty)$ ,  $E_{\text{est}}(\infty)$ ,  $E_{\text{l}}(\infty)$ ,  $a_{\text{u}}$ ,  $a_{\text{est}}$  and  $a_{\text{l}}$  are fitting parameters, and  $n_0$  is fixed to the average value of  $c_1$  in Eq. (3), with  $n_0 = -0.838$ . The two free parameters in each equation are obtained for each cluster from CCSD(T) energies calculated using the two largest available basis sets, resulting in an estimated CBS limit given by  $E_{\text{est}}(\infty)$  and an estimated error of  $[E_{\text{u}}(\infty) - E_{\text{l}}(\infty)]/2$ .

For the subset of small clusters with  $N = 2$ –8 the “efficient” procedure provides total energies that agree with the “accurate” values to within the estimated error ranges. The mean absolute errors for small clusters using “efficient” (TQ), (Q5) and (56) extrapolation are 6.1, 2.6, and 1.0 meV/[H<sub>2</sub>], respectively, where we use the notation  $(n_1 n_2)$  to denote “efficient extrapolation” from results with the cc-pVn<sub>1</sub>Z and cc-pVn<sub>2</sub>Z basis sets. The “efficiently” estimated energies and confidence intervals for each small cluster geometry are shown in Fig. 1 as a deviation from the “accurate” estimates.

The convergence of the DFT total energies with basis set is exponential due to the correlation being contained in the exchange-correlation functional rather than

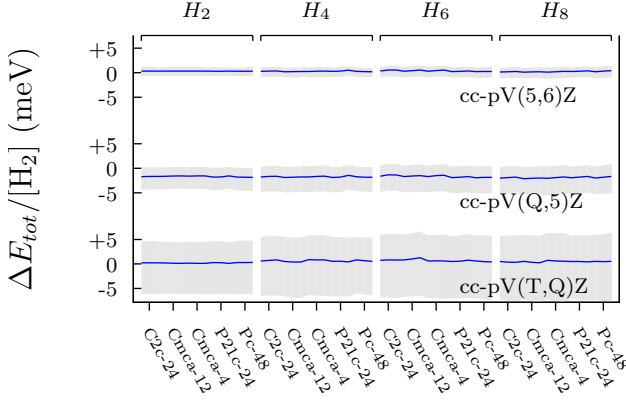


FIG. 1. “Efficient” estimates of CCSD(T) total energies, and the associated confidence intervals, plotted as the deviation from “accurate” total energies. Estimates are given for three different pairs of basis sets, and for geometries containing 8 hydrogen atoms or fewer. The gray shaded regions show the error ranges associated with the “efficient extrapolation” procedure.

the wave function itself. Consequently, extrapolation using only Eq. (2) provides accurate estimates of the CBS limit (replacing  $E_{\text{HF}}$  with the DFT energy). In a similar manner to the CCSD(T) calculations, we begin with an “accurate” estimate of the finite basis correction for the geometries with  $N = 2$ –12. The CBS total energy is provided by a least-squares fit of Eq. (2) to results obtained with the  $n = \text{D, T, Q, 5}$  basis sets, and with all parameters free.

The absolute deviation  $|E_{\text{KS}}(5) - E_{\text{KS}}(\infty)|$  is expected to severely overestimate the extrapolation error, hence mean and maximum absolute deviations of 1.4 and 4.4 meV/[H<sub>2</sub>] for clusters of 2–12 hydrogen atoms reliably validate this finite-basis correction. Note that these mean and maximum deviations are taken over all of the exchange-correlation functionals considered as well as the range of cluster sizes.

Then, we construct an “efficient” estimate of the CBS limit by two-point extrapolation using Eq. (2) with  $b_0$  fixed to its average value for the “accurate” extrapolations, where  $b_0 = 6.707$ . Upper and lower limits are provided by extrapolation using the maximum and minimum  $b_0$  values ( $b_0 = 8.247$  and  $4.257$ ). This provides estimated values and error ranges in a similar manner to the CCSD(T) approach, with mean errors of 19.7, 2.6, and 0.9 meV/[H<sub>2</sub>] for (DT), (TQ), and (Q5) extrapolations, respectively.

Further basis set errors are expected to remain after extrapolation due to superposition and the finite range of the cc-pVnZ basis sets. These are found to be negligible (see Supplementary Material<sup>35</sup>).

### III. OPTIMUM EXCHANGE-CORRELATION FUNCTIONALS

New exchange-correlation functionals are constructed by taking the components of the B3LYP functional<sup>36</sup> and determining their coefficients by minimizing the average difference between the CCSD(T) and DFT energies using an iterative approach. We use the B3LYP form as it is the most accurate of the standard hybrid functionals for hydrogen clusters, and parameterizations with and without exact HF exchange are well established. At iteration  $i$ , the total DFT energy provided by these B3LYP-like functionals is of the form

$$E_{\text{KS}}[\rho_i] = T[\rho_i] + V[\rho_i] + \alpha_0^{(i)} E_{\text{x}}^{\text{HF}}[\rho_i] + \alpha_1^{(i)} E_{\text{x}}^{\text{B88}}[\rho_i] + \alpha_2^{(i)} E_{\text{x}}^{\text{LDA}}[\rho_i] + \alpha_3^{(i)} E_{\text{c}}^{\text{LYP}}[\rho_i] + \alpha_4^{(i)} E_{\text{c}}^{\text{LDA}}[\rho_i] + E_{\text{D3}}, \quad (9)$$

where  $\rho_i$  is the electronic density,  $\alpha_j^{(i)}$  are linear parameters,  $E_{\text{x}}^{\text{HF}}$  is the exact HF exchange energy,  $E_{\text{x}}^{\text{B88}}$  and  $E_{\text{c}}^{\text{LYP}}$  are the BLYP-parameterized exchange and correlation functionals, and  $E_{\text{x}}^{\text{LDA}}$  and  $E_{\text{c}}^{\text{LDA}}$  are the parameterized local density approximation (LDA) exchange and correlation functionals. The final term is a dispersion correction,  $E_{\text{D3}}$ . It is usually desirable for the resulting functional to accurately reproduce the exchange and correlation energies of the homogeneous electron gas (HEG), which can be accomplished by enforcing the constraints  $\alpha_0 + \alpha_1 + \alpha_2 = 1$  and  $\alpha_3 + \alpha_4 = 1$ .

The iterative method works as follows. Parameter values at the start of iteration  $i$ ,  $\{\alpha_j^{(i)}\}$ , provide a DFT density  $\rho_i$ . We then define the penalty function,

$$P_i(\{\alpha_j^{(i)}\}) = \sum_k \left[ \frac{E_{\text{KS}}^k[\rho_i](\{\alpha_j^{(i)}\}) - E_{\text{CCSD(T)}}^k}{N_k} \right]^2, \quad (10)$$

where index  $k$  runs over all 180 clusters considered, and  $E_{\text{CCSD(T)}}^k$  is the best available “efficiently extrapolated” CCSD(T) energy for the  $k$ th cluster. “Efficient” (DT) extrapolation of DFT energies is used to limit the computational cost.

The penalty function is then minimized with respect to  $\{\alpha_j^{(i)}\}$  with the density fixed to  $\rho_i$ . Due to linearity this is easily achieved using matrix diagonalization, and provides the optimum fixed-density parameters for  $\rho_i$ ,  $\{\alpha_j^{(i+1)}\}$ . This process is then repeated until the variation in parameter values and total energies is sufficiently small.

We assess the effect of three different constraints on the quality of the functional, whose combination give rise to eight optimized functionals. Four of the eight optimum functionals generated (O<sub>1</sub>, O<sub>2</sub>, O<sub>3</sub>, and O<sub>4</sub>) do not include dispersion corrections ( $E_{\text{D3}} = 0$ ), while the rest (O<sub>1</sub>-D3, O<sub>2</sub>-D3, O<sub>3</sub>-D3, and O<sub>4</sub>-D3) include that of Grimme *et al.*<sup>37</sup>. Dispersion corrections optimized for

each functional are used when available, and when unavailable BLYP corrections are used for GGA functionals (LDA, PW91, O<sub>3</sub>-D3 and O<sub>4</sub>-D3) and B3LYP corrections for hybrid functionals (SOGGA11-X, B3H, M08-SO, O<sub>1</sub>-D3 and O<sub>2</sub>-D3).

We generate four hybrid functionals [O<sub>1</sub>(-D3) and O<sub>2</sub>(-D3)] that include exact exchange ( $\alpha_0$  optimizable) and four GGA functionals [O<sub>3</sub>(-D3) and O<sub>4</sub>(-D3)] that exclude exact exchange ( $\alpha_0 = 0$ ). Finally, four of the functionals [O<sub>1</sub>(-D3) and O<sub>3</sub>(-D3)] are constrained to obey the HEG limit ( $\alpha_2 = 1 - \alpha_0 - \alpha_1$  and  $\alpha_4 = 1 - \alpha_3$ ), and the remaining four [O<sub>2</sub>(-D3) and O<sub>4</sub>(-D3)] are not required to obey the HEG limit ( $\alpha_2$  and  $\alpha_4$  optimizable).

Convergence is rapid for all eight functionals, with changes in the total energy falling to less than  $10^{-3}$  meV/[H<sub>2</sub>] after four iterations. Optimum parameters and constraints are reported in Table I.

The error in the total energy estimates used for optimization is small. The CBS total energies obtained with (DT) extrapolation incur a maximum estimated error of 36 meV/[H<sub>2</sub>]. This is an overestimate; the maximum difference between energies obtained using (DT) extrapolation and those obtained using larger basis sets (see the next section) is 8 and 13 meV/[H<sub>2</sub>] for the optimum functionals that satisfy or do not satisfy the HEG limit, respectively.

We note that the parameters of the optimized functionals differ significantly from their B3LYP counterparts. For instance, the coefficient of the HF exchange term,  $\alpha_0$ , when present, is larger than for most hybrid functionals, which reflects the importance of self-interaction (SI) in hydrogen systems. The negative coefficients obtained when the HEG limit is not imposed are somewhat undesirable, but not physically significant, as the total exchange and correlation energies remain positive.

#### IV. RESULTS

We quantify the performance of exchange-correlation functionals by their ability to reproduce the CCSD(T) total energies for the H<sub>2</sub> clusters. In addition to our optimum functionals, we test 11 standard functionals readily available in the literature: one LDA, four GGA, and six hybrid functionals that combine exact exchange and semi-local exchange-correlation functionals. We test the functionals with and without dispersion corrections, providing a total of 30 distinct DFT Hamiltonians (see Table II).

Calculations are performed using cc-pVnZ basis sets with different  $n$  for different clusters, and for CCSD(T) and DFT. “Efficient” extrapolation to the CBS limit, together with the estimation of the extrapolation error, is performed as described previously.

In the CCSD(T) calculations we use (56), (Q5), and (TQ) extrapolations for clusters of  $N = 2-8$ ,  $10-12$ , and  $14-24$  atoms, respectively, obtaining mean basis set errors in the extrapolated total energies of 1.0, 2.9, and

7.3 meV/[H<sub>2</sub>]. In the DFT calculations we use (Q5) and (TQ) extrapolations for clusters of  $N = 2-12$  and  $14-24$  atoms, respectively, obtaining mean errors in the DFT energies of 0.9 and 2.6 meV/[H<sub>2</sub>]. (For DFT mean values of errors are evaluated over all functionals, as well as geometries.) This energy resolution is sufficient to assess the accuracy of exchange-correlation functionals.

Each error is quantified as the deviation of the DFT energy from the CCSD(T) energy,  $\Delta = E_{\text{KS}} - E_{\text{CCSD(T)}}$ . Errors for each functional are summarized as both the mean and maximum absolute difference for all clusters in Fig. 2 for functionals with and without dispersion corrections.

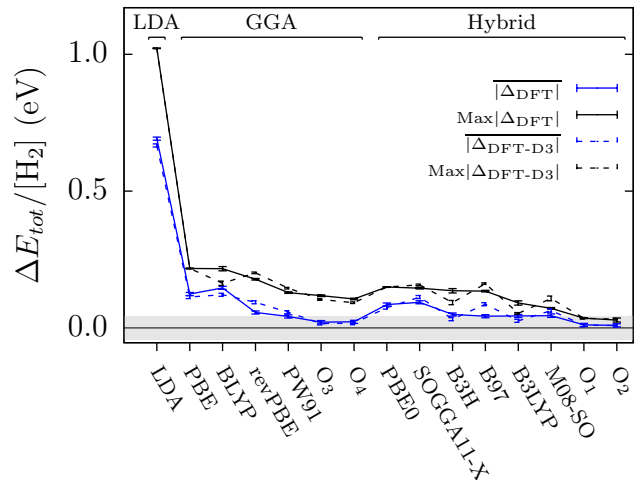


FIG. 2. Maximum and mean absolute differences between DFT and CCSD(T) total energies for a variety of exchange-correlation functionals, both including ( $\Delta_{\text{DFT-D3}}$ ) and excluding ( $\Delta_{\text{DFT}}$ ) dispersion corrections. The maximum and mean is taken over the 180 geometries considered, and basis set errors are almost imperceptible. Differences in the gray shaded region denote an accuracy of better than the chemical accuracy tolerance of  $1 \text{ kcal mol}^{-1} = 43 \text{ meV/[H}_2\text{]}$ .

As expected, LDA energies are significantly worse than those from all GGA and hybrid functionals, resulting in a peak error of  $\sim 1 \text{ eV/[H}_2\text{]}$  for H<sub>2</sub> that slowly falls with increasing cluster size. Most of the LDA error may be identified with an imperfect correction to the SI included in the Coulomb potential for DFT as this results in an error of  $1.2 \text{ eV/[H}_2\text{]}$  for an isolated hydrogen atom. The improved accuracy of the non-LDA functionals is, in part, due to their improved treatment of SI.

The general trend in Fig. 2 is for hybrid functionals to be more accurate than GGA functionals, with the optimized functionals providing the smallest errors for both types. However, there is considerable overlap between errors for the two functional types. Of the hybrid functionals, only B3LYP and the optimized functionals are more accurate than the optimized GGA functionals, whether or not dispersion is present.

Given the fitting procedure and the extra variational freedom allowed by relaxing the HEG limit, we would



Functional	Exchange			Correlation	
	$\alpha_0$	$\alpha_1$	$\alpha_2$	$\alpha_3$	$\alpha_4$
O <sub>1</sub>	0.60354910205605	0.14589659222789	$(1 - \alpha_0 - \alpha_1)$	0.54819918079350	$(1 - \alpha_3)$
O <sub>2</sub>	0.76679614121531	-0.84277505048128	1.05244207765216	-1.85005839047496	2.33023683529908
O <sub>3</sub>	(0)	0.79872851641853	$(1 - \alpha_0 - \alpha_1)$	0.61544187700225	$(1 - \alpha_3)$
O <sub>4</sub>	(0)	1.29167530141040	-0.25869459399816	2.03890740504165	-0.86103233660732
O <sub>1</sub> -D3	0.49200298266553	0.34311594882336	$(1 - \alpha_0 - \alpha_1)$	0.67727856886346	$(1 - \alpha_3)$
O <sub>2</sub> -D3	0.61183767165887	-0.39108315837275	0.76213024756349	-1.10597328426091	1.72092162085949
O <sub>3</sub> -D3	(0)	0.90916334640647	$(1 - \alpha_0 - \alpha_1)$	0.78790801474948	$(1 - \alpha_3)$
O <sub>4</sub> -D3	(0)	1.38553958317866	-0.35884244275049	2.16733823510143	-0.96208190626903
B3LYP	0.20000000000000	0.72000000000000	0.08000000000000	0.81000000000000	0.19000000000000

TABLE I. Parameters that define the optimized functionals used. Parameters for the standard B3LYP functional are also shown, for comparison.

Functional	Type
LDA	Local
PBE <sup>38</sup>	GGA
BLYP <sup>39</sup>	GGA
revPBE <sup>40</sup>	GGA
PW91 <sup>41</sup>	GGA
O <sub>3</sub> (-D3)	GGA
O <sub>4</sub> (-D3)	GGA
PBE0 <sup>42</sup>	Hybrid
SOGGA11-X <sup>43</sup>	Hybrid
B3H <sup>44</sup>	Hybrid
B97 <sup>45</sup>	Hybrid
B3LYP <sup>36</sup>	Hybrid
M08-SO <sup>46</sup>	Hybrid (meta)
O <sub>1</sub> (-D3)	Hybrid
O <sub>2</sub> (-D3)	Hybrid

TABLE II. Functionals used and their types.

expect O<sub>2</sub>-D3 and O<sub>4</sub>-D3 functionals to reproduce the CCSD(T) results to the highest accuracy. This is the case, with the lowest peak error for the hybrid functionals of 32(3) meV/[H<sub>2</sub>] provided by O<sub>2</sub>-D3, and the lowest peak error for the GGA functionals of 92(3) meV/[H<sub>2</sub>] provided by O<sub>4</sub>-D3. However, the O<sub>1</sub>-D3 and O<sub>3</sub>-D3 functionals, which are constrained to reproduce the HEG limit, perform only marginally worse, with peak errors of 35(3) meV/[H<sub>2</sub>] and 104(3) meV/[H<sub>2</sub>], respectively.

Dispersion provides no consistent improvement to the performance of the functionals considered. The introduction of dispersion increases the peak error for 5 of the 15 functionals, leaves it unchanged for another 5, and decreases it for the other 5 functionals. Unsurprisingly, functionals optimized with dispersion give marginally better results when applied with dispersion than without, and vice versa. This effect is not shown in the figures.

In order to gain further insight into the performance

of the functionals, in Fig. 3 we plot the deviation of the DFT energies using the O<sub>n</sub>(-D3) functionals from the CCSD(T) energies for each of the 180 clusters. Data generated with and without dispersion exhibit the same structure, confirming that dispersion has a small effect on the results. Overall, the hybrid functionals reproduce the CCSD(T) for all clusters to within chemical accuracy, typically associated with a tolerance of 1 kcal mol<sup>-1</sup> = 43 meV/[H<sub>2</sub>].

For the optimum GGA functionals, O<sub>3</sub>(-D3) and O<sub>4</sub>(-D3), the deviations from CCSD(T) are larger in magnitude, occasionally exceeding the chemical accuracy tolerance, and display more structure than those for the hybrid functionals. Figure 3 shows segments of steep (approximately) linear variation with the pressure of the underlying bulk structure, where each segment is associated with a different symmetry and cluster size. This dependence is barely apparent for the hybrid functionals.

The inaccuracy of the functionals considered is most apparent for cluster with 4–10 hydrogen atoms, and particularly for the H<sub>10</sub> clusters. This characteristic feature is identifiable for all of the GGA functionals considered, and for most of the hybrid functionals [it is not apparent for the B3H, M08-SO, O<sub>1</sub>(-D3), and O<sub>2</sub>(-D3) functionals].

For a number of the functionals [though not the O<sub>n</sub>(-D3)] there is an underlying linear variation of error with cluster size. This observation, together with the known inadequate description of SI by DFT, suggests that the SI error should be included in our assessment of the accuracy of the functionals.

We estimate the SI error as the difference between the total DFT energy of an isolated hydrogen atom evaluated with each functional (using an  $n = 6$  finite basis set that results in a negligible basis set error) and the exact analytic energy.

Figure 4 shows the resulting SI error for each functional (dispersion is zero for an isolated atom, so there are only 19 distinct total energies to consider). The greatest SI error arises for the LDA functional. For the optimum

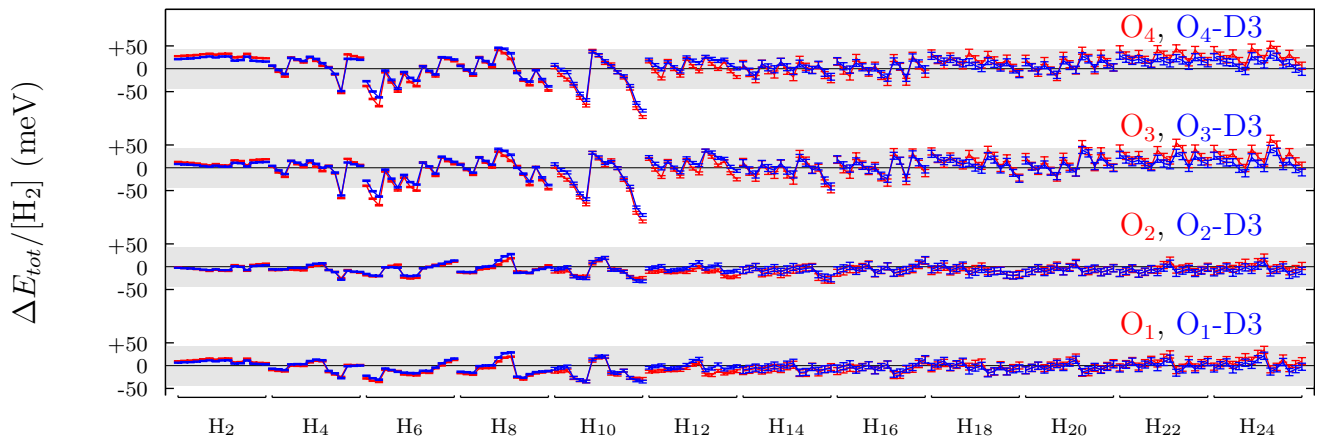


FIG. 3. Deviation of DFT total energies from the CCSD(T) total energies. Results with (blue) and without (red) dispersion are shown. Values and errors bars are obtained using the “efficient” estimates described in the text, and differences in the gray shaded regions denote an accuracy of better than the chemical accuracy tolerance of  $1 \text{ kcal mol}^{-1} = 43 \text{ meV}/[\text{H}_2]$ .

functionals the SI error is consistently decreased by including HF exchange and not enforcing the HEG limit. Of all the functionals considered the SI error is smallest for the  $\text{O}_2$  and  $\text{O}_2\text{-D3}$  hybrid functionals, taking values of 13 and 3  $\text{meV}/[\text{H}_2]$ , respectively.

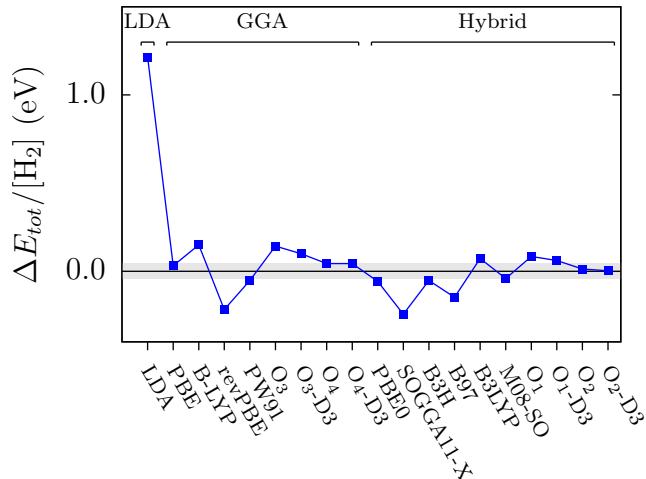


FIG. 4. Self-interaction errors for an isolated hydrogen atom for the 19 functionals considered in the text, and using the exact density. Points within the gray shaded region denote a self-interaction error smaller than the chemical accuracy tolerance of  $1 \text{ kcal mol}^{-1} = 43 \text{ meV}/[\text{H}_2]$ .

Since we are seeking an accurate functional for systems very unlike the HEG, relaxing the requirement that the HEG limit is conserved in the functionals is justified. However, it is reasonable to expect that a physically realistic functional will not radically misrepresent the HEG. Whether this is so may be assessed by summing the exchange and correlation coefficients in Table I. For all of the optimized functionals the HEG exchange is well preserved, with total HEG exchange deviating from the true value by less than 4%. The HEG correlation is

significantly less accurate when the HEG limit is not enforced, deviating from the true value by more than 17% for the GGA functionals, and more than 38% for the hybrid functionals. Consequently, we do not recommend using  $\text{O}_2$  or  $\text{O}_4$  despite the marginal improvement in accuracy they provide.

As a final test of the accuracy of the functionals we relax the geometry of the  $\text{H}_2$  molecule (using the cc-pV6Z basis set). All functionals except LDA reproduce CCSD(T) bond lengths to within chemical accuracy (of  $0.01 \text{ \AA}$ ). Errors in relaxed geometry molecular dissociation energies provided by the standard functionals fall in the range 77–206  $\text{meV}/[\text{H}_2]$  for GGA, and 10–143  $\text{meV}/[\text{H}_2]$  for hybrid functionals. For  $\text{O}_3$  and  $\text{O}_1$  the errors are 10 and 5  $\text{meV}/[\text{H}_2]$ , and halved by including dispersion. These results are beyond the energetics of fixed hydrogen clusters, and confirm that the improved accuracy can be expected to remain for other geometries.

The proposed functionals can be used to generate single-particle orbitals for quantum Monte Carlo calculations<sup>4</sup>, which is particularly useful for studies of crystalline hydrogen structures, such as those of Refs. 30 and 28, which require more accuracy than afforded by DFT and much larger systems sizes than are accessible with CCSD(T). In order to assess the accuracy of the functionals in this context we perform diffusion Monte Carlo (DMC) calculations using the CASINO code<sup>47</sup> for the  $\text{H}_6$ ,  $\text{H}_{12}$ , and  $\text{H}_{24}$  clusters using single-particle orbitals generated with the  $\text{O}_3$  GGA functional, as well as with the LDA, PBE, and B3LYP functionals, for comparison. The orbitals are cusp-corrected<sup>48</sup> to prevent divergences of the local energy at electron-nucleus coalescence points, and used in a Slater-Jastrow<sup>49,50</sup> trial wave function, whose parameters are optimized using linear least-squares energy minimization<sup>51,52</sup> at the variational Monte Carlo level. The DMC energies are obtained using the recent modifications of Zen *et al.*<sup>53</sup>, and extrapolated to zero timestep<sup>54</sup> and infinite population size; additional

details are given in the Supplementary Material<sup>35</sup>.

Our results show that the DMC energies are insensitive to the choice of DFT functional used to generate the single-particle orbitals to within statistical uncertainty. In Fig. 5 we plot the DMC energies corresponding to the  $O_3$  functional, which are higher than the CCSD(T) energies for all 45 clusters. Each of the successive S, D, and (T) contributions to the CCSD(T) energy is about an order of magnitude smaller than the previous one, which is expected in well-converged calculations, hence the missing correlation energy in our CCSD(T) results may reasonably be identified with the basis set error. Figure 5 also shows close agreement between DMC and CCSD energies, strongly suggesting that the small correlation energy absent from our DMC results can be identified with the triple excitations included in CCSD(T).

The energy differences between CCSD(T) and DMC are smaller than  $1 \text{ kcal mol}^{-1}$ , indicating that the single-particle orbitals generated with any of the tested DFT functionals, including our optimized  $O_3$  functional, are an excellent choice for performing accurate DMC calculations. The pressure dependence of the DMC energies can be attributed to the loss of accuracy of the DMC method at higher pressures.

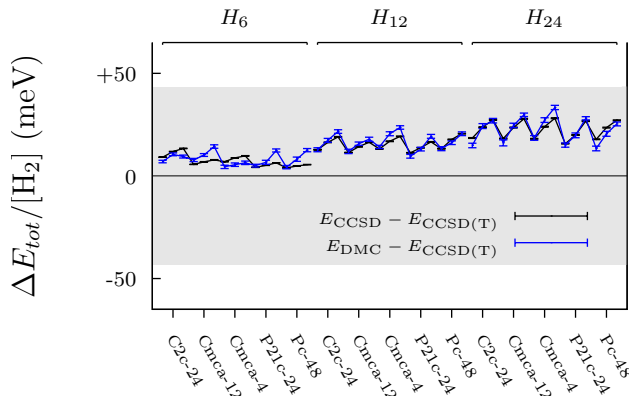


FIG. 5. DMC and CCSD energies relative to CCSD(T) for the  $H_6$ ,  $H_{12}$ , and  $H_{24}$  clusters using single-particle orbitals obtained with the  $O_3$  GGA functional. The error bars represent the statistical uncertainty in the DMC data, and the estimated basis-set extrapolation error in the CCSD data. All energy differences are smaller than the chemical accuracy tolerance of  $1 \text{ kcal mol}^{-1} = 43 \text{ meV}/[H_2]$  indicated by the gray shaded region.

In summary, the optimized functionals provide the best description of the total energies for the 180 clusters considered, with optimized hybrid functionals being the most accurate. Relaxing the natural requirement that the functionals reproduce the HEG limit provides an insignificant or small increase in the accuracy of the functionals, at the cost of an inaccurate description of the HEG correlation. Dispersion corrections are small for the optimized functionals.

These points suggest that the functionals of choice for hydrogen systems is the hybrid functional  $O_1(-D3)$ , with

$O_3(-D3)$  a second choice if a GGA is preferred. The two different forms are available for application with or without dispersion corrections. With dispersion excluded, these provide an accuracy of better than 118(3) and 36(1) meV/ $[H_2]$  for GGA and hybrid functionals, respectively. With dispersion included, these provide an accuracy of better than 104(3) and 35(3) meV/ $[H_2]$  for GGA and hybrid functionals, respectively.

The CCSD(T) energies of 18–24-atom clusters extracted from bulk *Cmca*-12 fall within 30 meV/ $[H_2]$  of each other, signaling convergence with system size. Also, while the standard density functionals considered in our work yield DFT errors in energy which steadily increase with system size, our optimized functionals do not exhibit this behavior. We conclude from these observations that the optimized functionals are likely to remain accurate in the bulk limit.

There are several obvious options to improve the accuracy of optimized functionals for hydrogen clusters. The most evident improvement would be to reduce the remnant of basis set error that arises from the extrapolation to the CBS limit used during functional optimization. This would be computationally expensive, and is probably not significant given that the structure shown in Fig. 4 is accurately reproduced by (DT), (TQ), or (Q5) extrapolation, so is not due to basis set error. Another simple option would be to dynamically allocate weights in the optimization process in order to equalize the distribution of errors between clusters. This would reduce the peak error at the cost of introducing a larger error for clusters containing 12 hydrogen atoms or more.

Perhaps the most promising option for improving the functionals would be to add more variational freedom to the optimized functional. The simplest approach would be to include other parameterizations of semi-local functionals, including meta terms, as this would preserve the linearity of the optimization we rely on within the iterative optimization process. A more general approach would be to directly optimize the enhancement factors present within the functional forms, introducing non-linearity into each optimization iteration. Similarly, the screening parameters present in the dispersion correction may also be optimized, but it seems likely that this will only provide a marginal improvement in the accuracy of the energies.

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